

Part II.

Applications of the method
of multipliers

What do we do next with the identity

$$\begin{aligned} & (\mathcal{D}_0 u, \mathcal{D}_0 u)_{L^2(M)} - (-\Delta_M u, u)_{L^2(M)} \\ &= \int_{\Omega} (2 \text{Jac}_{\vec{F}} [\nabla U, \nabla U] - |\nabla U|^2 \text{div} \vec{F}) \end{aligned} ?$$

For a fixed \vec{F} , the RHS is a quadratic form in $|\nabla U|$ with bounded coefficients, and thus its modulus is bounded above by $C \|\nabla U\|_{\Omega}^2 = C (\mathcal{D}_0 u, u)_M$

Thus we have

$$\left| (\mathcal{D}_0 u, \mathcal{D}_0 u)_M - (-\Delta_M u, u)_M \right| \leq C (\mathcal{D}_0 u, u)_M$$

C may be expressed via some geometric characteristics of Ω and M [Colbois/Gir. /Hass'18; Provenzano/Stubbe'19]

A functional analysis problem: given two non-negative self-adjoint operators A, B in a Hilbert space \mathcal{H}

$$\text{s.t. } |(Au, Au)_{\mathcal{H}} - (Bu, u)_{\mathcal{H}}| \leq C(Au, u)_{\mathcal{H}}$$

for all $u \in \text{Dom}(B) \cap \text{Dom}(A^2)$,

what can be said about the eigenvalues $\{\alpha_j\}$ of A and $\{\beta_j\}$ of B ?

Thm [Girouard/Karpukhin/ML/Polterovich '22] We have

$$|\alpha_k - \sqrt{\beta_k}| \leq C \quad \forall k \in \mathbb{N}.$$

Idea of proof: Use variational principles for both operators with test spaces constructed of eigenfunctions of the other operator

In our case, this implies that

$$|\hat{\sigma}_k(\Omega) - \sqrt{\lambda_k(-\Delta_M)}| \leq C \quad \forall k$$

Steklov e.v.

e.v. of the boundary Laplacian

Uniform bound! Recall that for the disk $C=0$

Remarks: • Smoothness assumptions on $M = \partial\Omega$

may be relaxed to $C^{2,\alpha}$, $\alpha > 0$ for $d > 2$
and $C^{1,1}$ for planar domains

- Analogues exist for domains on manifolds
- Will be important later for asymptotics of $\hat{\sigma}_k$ for large k .

We'll take a slight detour now to discuss the **parameter-dependent DtN map** \mathcal{D}_Λ .

Given $\Lambda \notin \text{Spec}(-\Delta_{\Omega}^{\text{Dir}})$ and $u \in H^{1/2}(\partial\Omega)$,

we can uniquely solve a non-homog. problem

$$\begin{cases} -\Delta U = \Lambda U & \text{in } \Omega \\ U|_M = u & \text{on } M = \partial\Omega \end{cases} \quad \left| \quad \begin{array}{l} U := \mathcal{E}_\Lambda u \text{ is called} \\ \Lambda\text{-Helmholtz extension} \\ \text{of } u \end{array} \right.$$

Then $\mathcal{D}_\Lambda: H^{1/2}(M) \rightarrow H^{-1/2}(M)$ is defined by

$$\mathcal{D}_\Lambda u = (\partial_n U)|_M = (\partial_n \mathcal{E}_\Lambda u)|_M$$

Main facts about \mathcal{D}_Λ :

- self-adjoint, with discrete real spectrum

$$\sigma_{\Lambda,1} \leq \sigma_{\Lambda,2} \leq \dots \leq \sigma_{\Lambda,k} \leq \dots \rightarrow +\infty$$

- $(\mathcal{D}_\Lambda u, u)_{L^2(M)} = \|\nabla U\|_{L^2(\Omega)}^2 - \Lambda \|U\|_{L^2(\Omega)}^2$

- variational principle

$$\sigma_{\Lambda,k} = \min_{\substack{\tilde{\mathcal{L}} \subset CH^{1/2}(M) \\ \dim \tilde{\mathcal{L}} = k}} \max_{\substack{u \in \tilde{\mathcal{L}} \\ u \neq 0}} \frac{(\mathcal{D}_\Lambda u, u)_{L^2(M)}}{\|u\|_{L^2(M)}^2}$$

$$= \min_{\substack{\mathcal{L} \subset \mathcal{H}_\Lambda(\Omega) \\ \dim \mathcal{L} = k}} \max_{\substack{U \in \mathcal{L} \\ U \neq 0}} \frac{\|\nabla U\|_\Omega^2 - \Lambda \|U\|_\Omega^2}{\|U\|_M^2}$$

$$\mathcal{H}_\Lambda(\Omega) := \{U \in H^1(\Omega) : -\Delta U - \Lambda U = 0\}$$

can be repl.
here by

$$H^1(\Omega) \text{ if } \Lambda < \lambda_1^D(\Omega)$$

- definition of \mathcal{D}_Λ can be extended to $\Lambda \in \text{Spec}(-\Delta_{\Omega}^{\text{Dir}})$ if we restrict its domain to the orthogonal complement of the subspace of normal derivatives of the corresponding Dirichlet eigenspace

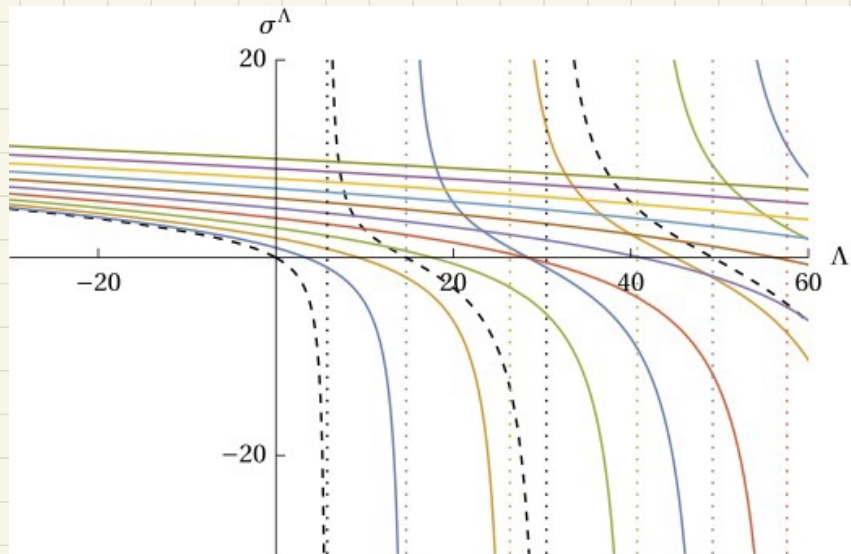
- DTN - Robin duality:

$$\mathcal{G} \in \text{Spec}(\mathcal{D}_\Lambda) \iff \Lambda \in \text{Spec}(-\Delta_{\Omega}^{\text{Robin}, -\mathcal{G}})$$

with the same multiplicities!

[Friedlander '91; Arendt/Mazzeo '12; Hassannezhad/Sher '22
+ history]

Example: Eigenvalues of \mathcal{D}_Λ for the unit disk as functions of Λ



Observations

- eigenvalue curves cross $\sigma = 0$ when $\Lambda \in \text{Spec}(-\Delta^{\text{Neu}})$
- eigenvalue curves "blow up" $\Lambda \in \text{Spec}(-\Delta^{\text{Dir}})$
- eigenvalue

Exercise: show that in this case the spectrum of \mathcal{D}_Λ consists of eigenvalues

$$\begin{cases} I'_m(-\sqrt{\Lambda}) / I_m(-\sqrt{\Lambda}), & \Lambda < 0 \\ m, & \Lambda = 0 \\ J'_m(\sqrt{\Lambda}) / J_m(\sqrt{\Lambda}), & \Lambda > 0 \end{cases}$$

single for $m = 0$
double for $m > 0$

from $-\infty$ to $+\infty$ when are monotone decreasing in Λ otherwise

These observations can be turned into rigorous thms not only for a disk but for a general Lipschitz Euclidean domain or a Riem. manifold with boundary, and in particular imply

Thm. [Friedlander; Arendt-Mazzeo]

$$W_{-2}^{\text{Neu}}(\Lambda) - W_{-2}^{\text{Dir}}(\Lambda) = W^{\mathcal{D}_\Lambda}(0)$$

number of negative eigenvalues of \mathcal{D}_Λ $\forall \Lambda \in \mathbb{R}$

Corollary. If $\Omega \subset \mathbb{R}^d$, then $\lambda_{k+1}^{\text{Neu}} < \lambda_k^{\text{Dir}} \quad \forall k \in \mathbb{N}$

Proof uses the fact that $W^{\mathcal{D}_\Lambda}(0) \geq 1$ for $\Lambda > \lambda_1^{\text{Dir}}$

(alternative elementary proof by [Filonov'04])

(the corollary may not hold in Riemannian case)

We now ask the following question:

Can we compare the eigenvalues of $\mathcal{D}_\Lambda(\Omega)$ with those of $-\Delta_{\partial\Omega}$ for (some) $\Lambda \neq 0$?

We have

Thm Let $\Omega \subset \mathbb{R}^d$ be a bdd domain with smooth bdry $M = \partial\Omega$. Then for $\Lambda \leq 0$ we have

$$|\sigma_k^\Lambda - \sqrt{\lambda_k(-\Delta_M) - \Lambda}| \leq C \quad k \in \mathbb{N}$$

with some constant C uniformly in
both k and Λ

Ideas of proof: [GKLP'22]

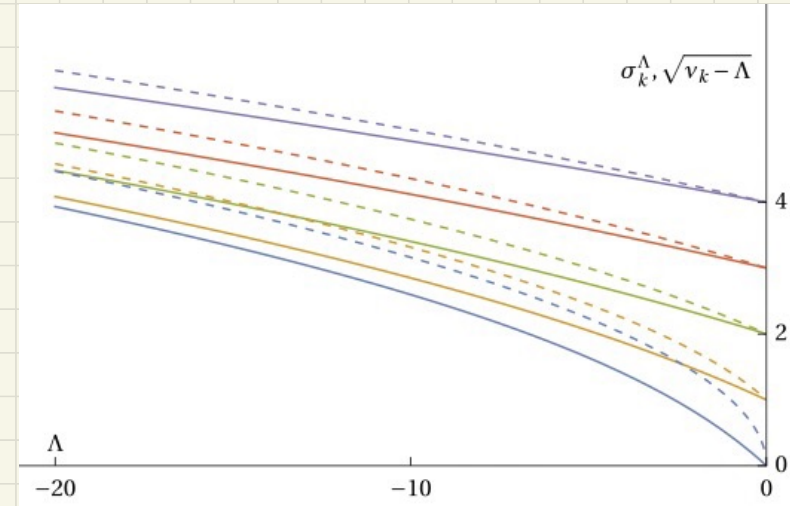
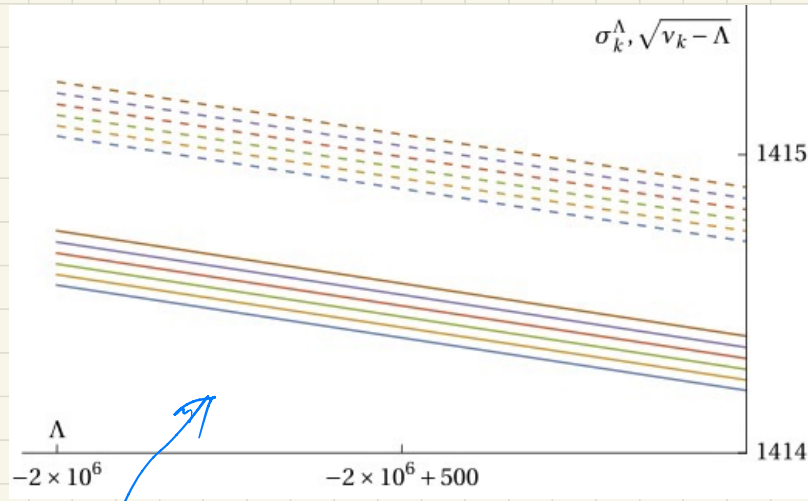
- A variant of generalised Pohožaev's identity for Helmholtz [Hassannezhad Siffert'20]
- Hence a variant of generalised Hörmander's inequality

$$\left| (\mathcal{D}_\lambda u, \mathcal{D}_\lambda u)_M - ((-\Delta_M - \lambda)u, u)_M \right| \leq C (\mathcal{D}_\lambda u, u)_M$$

- Use abstract bound with $\mathcal{A} = \mathcal{D}_\lambda$, $\mathcal{B} = -\Delta_M - \lambda$.

Remark We will see later that **no analogue** of this result may hold if boundary has corners

Illustration: the unit disk



solid curves: σ_k^Λ

$k \in \{100, 102, 104, 106, 108\}$

dashed curves: $\sqrt{\lambda_k^{-\Delta_M} - \Lambda}$

$k \in \{1, 3, 5, 7, 9\}$

Before we proceed, some further references (full bibliography will appear at the end of the last set of slides)

[Chandler-Wilde/Graham/Langdon/Spence '12] for a historical overview of the method of multipliers

[Hassell Tao '02] for applications to bounds of $\| \partial_n u_j^{\text{Dir}} \|_{L^2(\partial\Omega)}^2$ on normal derivatives of Dirichlet e.f.s

[Rudnick Wigman Yesha '21] for bounds on $\| u_j^{\text{Rob}, \gamma} \|_{L^2(\partial\Omega)}^2$ on traces of Robin e.f.s.

etc, etc, ...

Part III.

Spectral asymptotics

for the Dirac map \mathcal{D}_0

For the rest of this course we will be looking at the asymptotic behaviour of eigenvalues of the Steklov problem on $\Omega \subset \mathbb{R}^d$ (the DtN map \mathcal{D}_0), both in terms of asymptotics of eigenvalues

[mostly $d=2$]

σ_k , $k \nearrow +\infty$ and the counting function $\mathcal{N}^S(\sigma) := \{k : \sigma_k \leq \sigma\}$ as $\sigma \nearrow +\infty$

I'll start by listing some relatively well-known facts:

• Suppose that Ω has a smooth boundary M .

Then \mathcal{D}_0 is an **elliptic pseudodifferential operator of order 1**. Its **principal symbol** is given by $|\xi|$ and coincides with that of $\sqrt{-\Delta_M}$. Hence these two operators have the same leading term Weyl's asymptotics, and

$$\mathcal{N}^S(\mathcal{D}) = \underset{\substack{\uparrow \\ \text{Weyl constant}}}{C_{d-1}} |M|_{d-1} \mathcal{O}^{d-1} + o(\mathcal{O}^{d-1})$$

as $\mathcal{O} \rightarrow +\infty$

- By our previous comparison result, the same asymptotics holds if $\partial\Omega$ is not C^∞ but smooth enough
- On the other hand, for C^∞ boundary in the ~~smooth~~^{planar} case, error term is much better [Rozenblyum '86] [Edward 93 after Guillemin/Melrose]

$$G_k(\Omega) = G_k(\Omega^*) + o(k^{-N}) \quad \forall N$$

$k \rightarrow +\infty$

Ω^* — disk with the same perimeter as Ω

• For general Lipschitz domains in $d \geq 3$ the one-term Weyl's asymptotics is still an open conjecture.

In $d=2$ it was proved very recently by different techniques

[Karpukhin Lagacé Polterovich '22]

So the question is: what is happening for "not so smooth" planar domains, say polygons, and can we improve Weyl's asymptotics there?

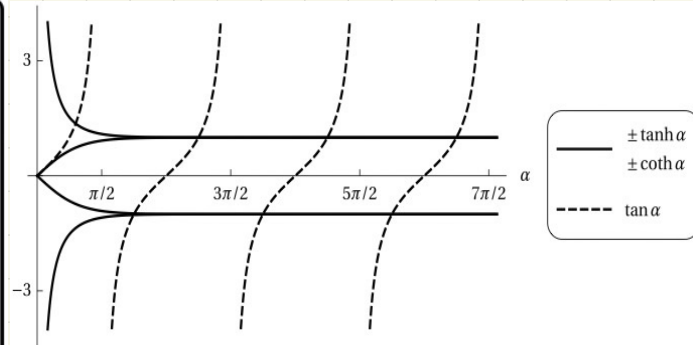
We start with a seemingly simple

Example. $\Omega = (-1, 1)^2$ a square.

[Girouard/Polterovich'17]

We try to find eigenfunctions by separation of variables which gives us

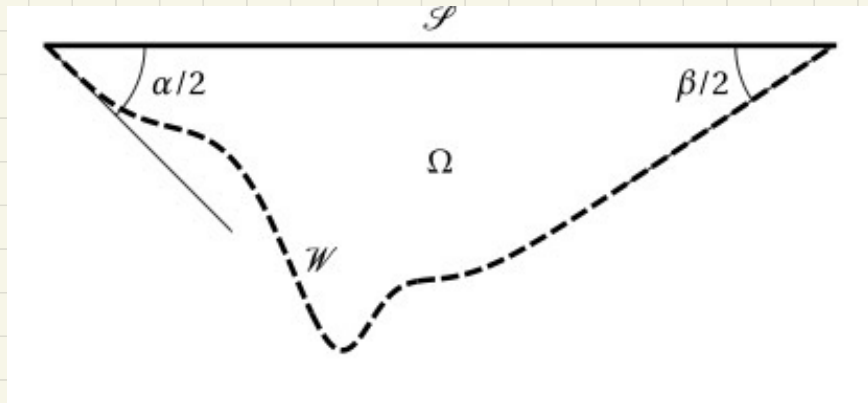
Eigenfunction	Equation for κ	Eigenvalue σ	Multiplicity
$U^0 := 1$		0	1
$U^1 := xy$		1	1
$U_\kappa^2 := \cos(\kappa x) \cosh(\kappa y)$	$\tan \kappa + \tanh \kappa = 0$	$\kappa \tanh \kappa$	2
$U_\kappa^3 := \cosh(\kappa x) \cos(\kappa y)$	$\tan \kappa - \coth \kappa = 0$	$\kappa \coth \kappa$	2
$U_\kappa^4 := \sin(\kappa x) \cosh(\kappa y)$	$\tan \kappa + \tanh \kappa = 0$	$\kappa \tanh \kappa$	2
$U_\kappa^5 := \cosh(\kappa x) \sin(\kappa y)$	$\tan \kappa - \coth \kappa = 0$	$\kappa \coth \kappa$	2
$U_\kappa^6 := \cos(\kappa x) \sinh(\kappa y)$	$\tan \kappa + \tanh \kappa = 0$	$\kappa \tanh \kappa$	2
$U_\kappa^7 := \sinh(\kappa x) \cos(\kappa y)$	$\tan \kappa - \coth \kappa = 0$	$\kappa \coth \kappa$	2
$U_\kappa^8 := \sin(\kappa x) \sinh(\kappa y)$	$\tan \kappa + \tanh \kappa = 0$	$\kappa \tanh \kappa$	2
$U_\kappa^9 := \sinh(\kappa x) \sin(\kappa y)$	$\tan \kappa - \coth \kappa = 0$	$\kappa \coth \kappa$	2



each intersection of a dotted curve with a solid curve gives a double e.v.

But how do we prove that we have found all the eigenvalues and haven't missed any?
 To do this, we have to take two sidesteps.

Slushing problem



$$\begin{cases} \Delta U = 0 & \text{in } \Omega \\ \partial_n U = g & \text{on } \mathcal{P} \\ \partial_n U = 0 & \text{on } \mathcal{W} \end{cases}$$

instead of pure Neumann on \mathcal{W} we may decompose

$$\mathcal{W}^p = \mathcal{W}_N \cup \mathcal{W}_D$$

19th century hydrodynamics!

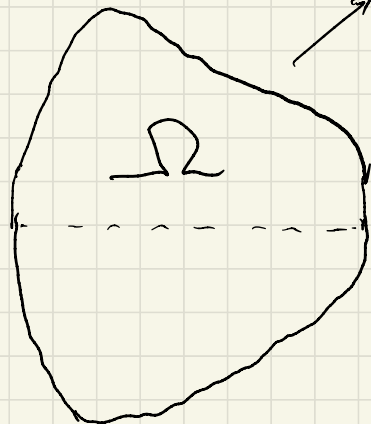
and impose Neumann here Dirichlet

Whatever reasonable b.c. we impose on \mathcal{W} , we always have

- the spectrum of the sloshing problem (or another mixed Steklov-Dirichlet-Neumann problem) is **discrete** and **non-negative**
- the eigenfunctions restricted to \mathcal{S} form a **basis** in $L^2(\mathcal{S})$

Sloshing problem will re-appear later!

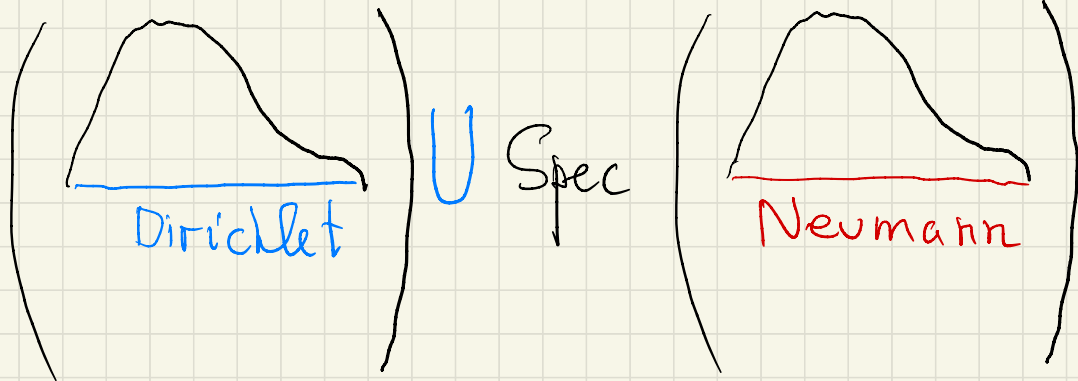
The second trick we need is the **symmetry reduction**



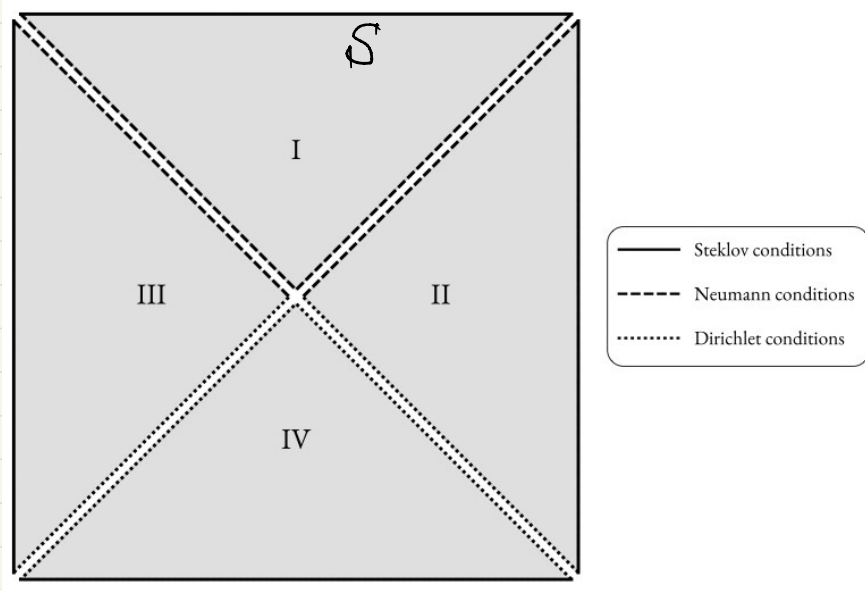
domain with a hyperplane of symmetry, problem and b.c. symmetric

then every eigenfunction is either **antisymmetric** or **symmetric**

Spec of Ω = Spec



Returning now to our Steklov problem on a square: using diagonals as lines of symmetry, we get



four mixed Steklov-Dir-Neumann problems in $45^\circ-45^\circ-90^\circ$ triangles.

Look at problem I:

from found eigen functions of the square we construct e.f.s $U^0, U^1, U^2_{\mathbb{R}} + U^3_{\mathbb{R}}, U^4_{\mathbb{R}} + U^5_{\mathbb{R}}$ for it.

Why is it the full set? Their traces on S form the full set of e.f. of $f^{(IV)}(x) = \mathbb{R}^4 f(x), f''(\pm 1) = f'''(\pm 1) = 0$ and therefore the basis in $L^2(S)$. Now repeat for II, III, IV...

So, we know that all eigenvalues of Steklov on $[-1, 1]^2$ are given by

Eigenfunction	Equation for κ	Eigenvalue σ	Multiplicity
$U^0 := 1$		0	1
$U^1 := xy$		1	1
$U_\kappa^2 := \cos(\kappa x) \cosh(\kappa y)$	$\tan \kappa + \tanh \kappa = 0$	$\kappa \tanh \kappa$	2
$U_\kappa^3 := \cosh(\kappa x) \cos(\kappa y)$			
$U_\kappa^4 := \sin(\kappa x) \cosh(\kappa y)$	$\tan \kappa - \coth \kappa = 0$	$\kappa \tanh \kappa$	2
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$U_\kappa^7 := \sinh(\kappa x) \cos(\kappa y)$			
$U_\kappa^8 := \sin(\kappa x) \sinh(\kappa y)$	$\tan \kappa - \tanh \kappa = 0$	$\kappa \coth \kappa$	2
$U_\kappa^9 := \sinh(\kappa x) \sin(\kappa y)$			

Corollary Steklov eigenvalues of this square satisfy

$$\sigma_{4m-k} = \left(m - \frac{1}{2}\right) \frac{\pi}{2} + O(m^{-\infty})$$

$m \in \mathbb{N}, k \in \{0, \dots, 3\}$ as $m \nearrow \infty$

Eigenvalues asymptotically come in clusters of 4

Q: Is it because we have 4 (equal) sides?

Would they appear in clusters of 5 for a regular pentagon?

Part IV.

Asymptotics of Steklov eigenvalues in curvilinear polygons

Most of the material in this part is covered by two long papers

ML + Parnowski + Polterovich + Sher

J. d'Anal. Math.

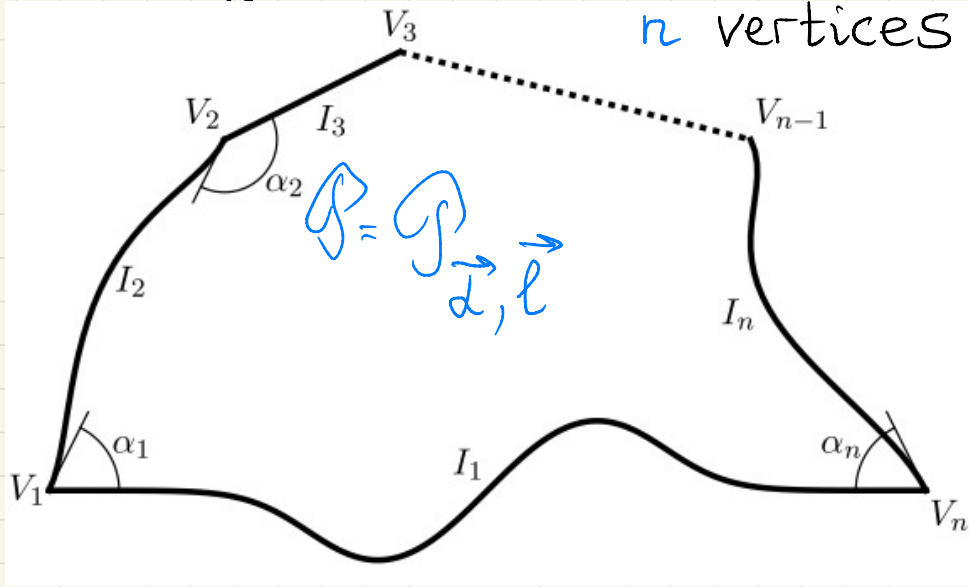
2021

Proc LMS

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but notation here is slightly different

Setting:



$$\left\{ \begin{array}{l} \Delta U = 0 \quad \text{in } \mathcal{G} \\ \partial_n U = \partial U \quad \text{on } \partial \mathcal{G} \end{array} \right.$$

curvilinear polygon

$$\mathcal{G}(\vec{\alpha}, \vec{l})$$

vector of angles

$$\vec{\alpha} = (\alpha_1, \dots, \alpha_n),$$

$$0 < \alpha_j < \pi$$

vector of sidelengths

$$\vec{l} = (l_1, \dots, l_n)$$

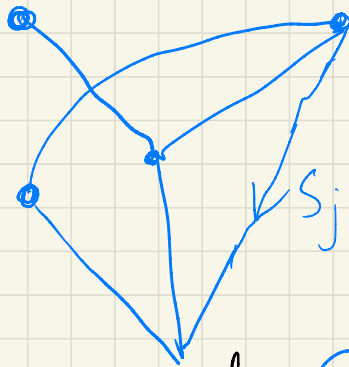
side I_j of length

l_j joins vertices

V_{j-1} and V_j

$$L = |\partial \mathcal{G}| = l_1 + \dots + l_n$$

Before stating the main result on the asymptotics of $G_m(\mathcal{G}_{\vec{\alpha}, \vec{\ell}})$, $m \rightarrow +\infty$ I need another side-step into the theory of **quantum graphs** [Berkolaiko & Kuchment '13]



metric graph G

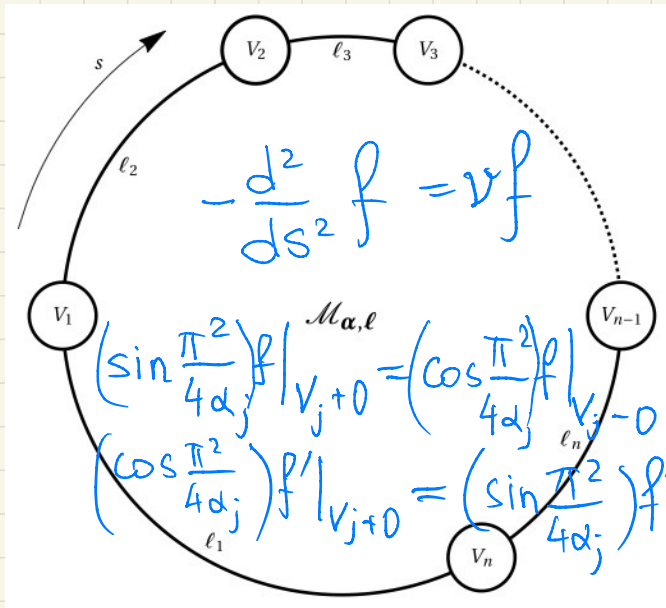
↑
edges E_j have lengths

$$-\Delta_G := \bigoplus \left(-\frac{d^2}{ds_j^2} \right) \Big|_{E_j}$$

plus some
self-adjoint boundary/
matching conditions

Looks deceptively simple —
in fact, deep subject!

Our main result, philosophy: Once more, compare the Steklov eigenvalues to those of a "boundary operator" but this time the boundary operator is a quantum graph associated with a curvilinear polygon $\mathcal{G}_{\vec{\alpha}, \vec{\ell}}$



Thm Let $\mathcal{G}_{\vec{\alpha}, \vec{\ell}}$ be a curvilinear polygon, σ_m its Steklov e.v.s and ν_m - e.v.s of QG $\mathcal{M}_{\vec{\alpha}, \vec{\ell}}$. Then

$$\sigma_m = \sqrt{\nu_m} + O(m^{-\epsilon}) \quad m \rightarrow +\infty$$

for some $\epsilon > 0$

As the QG $M_{\vec{\alpha}, \vec{\beta}}$ depends only on $\vec{\alpha}$ and $\vec{\beta}$, we have

Corollary If Ω^I, Ω^II are two curvilinear polygons with the same angles and sidelengths taken in the same order then

$$|G_m^I - G_m^{II}| \approx O(m^{-\varepsilon}) \quad m \nearrow +\infty$$

Defn From now on, the numbers $\tau_m := \sqrt{\lambda_m}$ are called the quasi-eigenvalues of the Steklov problem on G .