

Geometric bounds for Steklov eigenvalues

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We consider the Steklov eigenvalue problem on $\Omega \subset \mathbb{R}^N$

$$\begin{cases} \Delta u = 0, & \text{in } \Omega, \\ \frac{\partial u}{\partial \nu} = \sigma u, & \text{on } \partial\Omega. \end{cases}$$

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A. Girouard, I. Polterovich, *Spectral geometry of the Steklov problem*. Journal of Spectral Theory, 7 (2017), 321–359.

Variational characterization of Steklov eigenvalues:

$$\sigma_j = \min_{\substack{V \leq H^1(\Omega), \\ \dim V = j+1}} \max_{0 \neq u \in V} \frac{\int_{\Omega} |\nabla u|^2 dx}{\int_{\partial\Omega} u^2 d\sigma}$$

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Weyl's asymptotic law (if Ω is piecewise C^1):

$$\sigma_j \sim 2\pi\omega_{N-1}^{-\frac{1}{N-1}} \left(\frac{j}{|\partial\Omega|} \right)^{\frac{1}{N-1}} \text{ as } j \rightarrow +\infty,$$

where ω_{N-1} is the volume of the unit ball in \mathbb{R}^{N-1} .

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Remark: analogous inequalities hold for Dirichlet eigenvalues (Li-Yau, lower bounds), Neumann eigenvalues (Kröger, upper bounds) and eigenvalues of the Laplacian on Riemannian manifolds (Buser, Cheng-Yang, Colbois-Maerten).

The problem is completely solved for **simply connected** $\Omega \subset \mathbb{R}^2$:

$$\sigma_1 \leq \frac{2\pi}{|\partial\Omega|}$$

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In higher dimension we have an **isoperimetric control** of the eigenvalues

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where

$$I(\Omega) = \frac{|\partial\Omega|}{|\Omega|^{\frac{N-1}{N}}}$$

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We also have **isodiametric control**

$$\sigma_j \leq C_N \frac{j^{\frac{2}{N}+1}}{\text{diam}(\Omega)}$$

B. Bogosel, D. Bucur, A. Giacomini, *Optimal Shapes Maximizing the Steklov Eigenvalues*. SIAM J. Math. Anal., 49(2) (2017), 1645–1680.

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with the variational characterization

$$\lambda_j = \min_{\substack{V \subseteq H^1(\partial\Omega), \\ \dim V = j+1}} \max_{0 \neq u \in V} \frac{\int_{\partial\Omega} |\nabla_{\partial\Omega} u|^2 d\sigma}{\int_{\partial\Omega} u^2 d\sigma}$$

The eigenvalues satisfy the Weyl's asymptotic law

$$\lambda_j \sim 4\pi^2 \omega_{N-1}^{-\frac{2}{N-1}} \left(\frac{j}{|\partial\Omega|} \right)^{\frac{2}{N-1}} \text{ as } j \rightarrow +\infty.$$

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and Weyl-type bounds of the form

$$\lambda_j \leq \frac{(N-2)\kappa^2}{4} + C_N \left(\frac{j}{|\partial\Omega|} \right)^{\frac{2}{N-1}}.$$

P. Buser, *Beispiele für λ_1 auf kompakten Mannigfaltigkeiten*. Math. Z., 165(2) (1979), 107–133.

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Asymptotic formulas suggest that for large j

$$\sigma_j \approx \sqrt{\lambda_j}$$

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Theorem (P. - Stubbe, 2017)

Let $\Omega \subset \mathbb{R}^N$ be a bounded domain with connected boundary $\partial\Omega$ of class C^2 . Then there exists a constant c_Ω such that for all $j \in \mathbb{N}$

$$\lambda_j \leq \sigma_j^2 + 2c_\Omega \sigma_j, \quad \sigma_j \leq c_\Omega + \sqrt{c_\Omega^2 + \lambda_j}.$$

In particular,

$$|\sigma_j - \sqrt{\lambda_j}| \leq 2c_\Omega,$$

the constant c_Ω depending on the maximal possible size of a tubular neighborhood about $\partial\Omega$ and on the mean of the maximal curvatures of $\partial\Omega$.

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L. Provenzano, J. Stubbe, *Weyl-type bounds for Steklov eigenvalues*. arXiv:1611.00929 (2017).
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$$\bar{H}_\infty = \max_{x \in \partial\Omega} \left(\frac{1}{N-1} \sum_{i=1}^{N-1} |\kappa_i(x)| \right)$$

and $\kappa_i(x)$, $i = 1, \dots, N-1$ are the principal curvatures of $\partial\Omega$ at x .

Remark

One side of the estimate can be refined so that

$$\sigma_j - \sqrt{\lambda_j} \leq \frac{1}{h} + (N - 1)H_\infty$$

where $H_\infty := \max_{x \in \partial\Omega} |H(x)|$ with $H(x)$ the mean curvature of $\partial\Omega$ at x .

As a consequence of our main result we have a number of corollaries

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Corollary (Weyl-type upper bounds for Steklov eigenvalues)

Let Ω be a bounded domain of class C^2 in \mathbb{R}^N with connected boundary. Then for all $j \in \mathbb{N}$ it holds

$$\sigma_j \leq A_\Omega + C_N \left(\frac{j}{|\partial\Omega|} \right)^{\frac{1}{N-1}},$$

where $A_\Omega > 0$ depends on Ω and $C_N > 0$ depends only on the dimension N .

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The corollary follows from the main result and from upper bounds for Laplacian eigenvalues on $\partial\Omega$.

Theorem (Asymptotically sharp upper bounds for Steklov Riesz-means)

Let $\Omega \subset \mathbb{R}^N$ be a bounded domain with connected boundary $\partial\Omega$ of class C^2 . Then for all $z \geq 0$

$$\sum_{j=0}^{\infty} (z - \sigma_j)_+^2 \leq \frac{2}{N(N+1)} (2\pi)^{-(N-1)} \omega_{N-1} |\partial\Omega| (z + c_\Omega)^{N+1}.$$

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Bounds are asymptotically sharp since

$$\lim_{z \rightarrow +\infty} \frac{1}{z^{N+1}} \sum_{j=0}^{\infty} (z - \sigma_j)_+^2 = \frac{2}{N(N+1)} (2\pi)^{-(N-1)} \omega_{N-1} |\partial\Omega|.$$

Corollary (Sharp upper bounds for the trace of the Steklov heat kernel)

Let Ω be a bounded domain of class C^2 in \mathbb{R}^N with connected boundary.
Then

$$\sum_{j=0}^{\infty} e^{-\sigma_j t} \leq \frac{1}{N(N+1)} (2\pi)^{-(N-1)} \omega_{N-1} |\partial\Omega| t^{-N-1} e^{c_\Omega t} \Gamma(N+2, c_\Omega t)$$

for all $t > 0$, where $\Gamma(a, b) = \int_b^\infty t^{a-1} e^{-t} dt$.

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The estimate is sharp as $t \rightarrow 0^+$ since it implies

$$\limsup_{t \rightarrow 0^+} t^{N-1} \sum_{j=0}^{\infty} e^{-\sigma_j t} \leq (2\pi)^{-N-1} B_{N-1} \Gamma(N) |\partial\Omega|.$$

Corollary (Weyl-type lower bounds)

Let Ω be a bounded domain of class C^2 in \mathbb{R}^N with connected boundary. Then for all $j \in \mathbb{N}$:

$$\sigma_j \geq r_N 2\pi \omega_{N-1}^{-\frac{1}{N-1}} \left(\frac{j+1}{|\partial\Omega|} \right)^{\frac{1}{N-1}} - c_\Omega$$

with $r_N = \frac{N}{e\Gamma(N+1)^{1/N}} \leq 1$.

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This corollary is an immediate consequence of the bounds on the Steklov heat kernel.

The key point is to prove that for an **harmonic function** v in Ω the $L^2(\partial\Omega)$ -norm of the **normal derivative** and the **tangential gradient** are **equivalent**

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$$i) \quad \int_{\partial\Omega} |\nabla_{\partial\Omega} v|^2 d\sigma \leq \int_{\partial\Omega} \left(\frac{\partial v}{\partial \nu}\right)^2 d\sigma + 2c_{\Omega} \left(\int_{\partial\Omega} \left(\frac{\partial v}{\partial \nu}\right)^2 d\sigma\right)^{\frac{1}{2}}$$

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$$ii) \quad \left(\int_{\partial\Omega} \left(\frac{\partial v}{\partial \nu}\right)^2 d\sigma\right)^{\frac{1}{2}} \leq c_\Omega + \sqrt{c_\Omega^2 + \int_{\partial\Omega} |\nabla_{\partial\Omega} v|^2 d\sigma}$$

If we formally substitute $\lambda_j \sim \int_{\partial\Omega} |\nabla_{\partial\Omega} v|^2 d\sigma$ and $\sigma_j \sim \left(\int_{\partial\Omega} \left(\frac{\partial v}{\partial \nu}\right)^2 d\sigma\right)^{\frac{1}{2}}$,
i) and *ii)* become

$$\lambda_j \leq \sigma_j^2 + 2c_\Omega \sigma_j, \quad \sigma_j \leq c_\Omega + \sqrt{c_\Omega^2 + \lambda_j}.$$

The bounds follow from the min-max principle for σ_j and λ_j and the fact that both the Laplace-Beltrami eigenfunctions and the restrictions of the Steklov eigenfunction to $\partial\Omega$ form a Hilbert basis of $L^2(\partial\Omega)$.

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In particular, for λ_j :

$$\lambda_j = \inf_{\substack{V \leq H^1(\partial\Omega) \\ \dim V = j+1}} \sup_{\substack{0 \neq v \in V \\ \int_{\partial\Omega} v^2 d\sigma = 1}} \int_{\partial\Omega} |\nabla_{\partial\Omega} v|^2 d\sigma$$

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where $V_S = \langle u_0|_{\partial\Omega}, \dots, u_j|_{\partial\Omega} \rangle$, u_i are the first $j+1$ Steklov eigenfunctions on $\partial\Omega$.

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where $V_{LB} = \langle \phi_0, \dots, \phi_j \rangle$, ϕ_i are the harmonic extension in Ω of the first $j+1$ Laplace-Beltrami eigenfunctions on $\partial\Omega$. Since $v \in V_{LB}$ are harmonic,

$$\int_{\Omega} |\nabla v|^2 dx \leq \left(\int_{\partial\Omega} \left(\frac{\partial v}{\partial \nu} \right)^2 d\sigma \right)^{\frac{1}{2}}$$

and the bounds follow from *ii*).

For σ_j :

$$\sigma_j = \inf_{\substack{V \leq H^1(\Omega) \\ \dim V = j+1}} \sup_{\substack{0 \neq v \in V \\ \int_{\partial\Omega} v^2 d\sigma = 1}} \int_{\Omega} |\nabla v|^2 dx \leq \sup_{\substack{0 \neq v \in V_{LB} \\ \int_{\partial\Omega} v^2 d\sigma = 1}} \int_{\Omega} |\nabla v|^2 dx$$

where $V_{LB} = \langle \phi_0, \dots, \phi_j \rangle$, ϕ_i are the harmonic extension in Ω of the first $j+1$ Laplace-Beltrami eigenfunctions on $\partial\Omega$. Since $v \in V_{LB}$ are harmonic,

$$\int_{\Omega} |\nabla v|^2 dx \leq \left(\int_{\partial\Omega} \left(\frac{\partial v}{\partial \nu} \right)^2 d\sigma \right)^{\frac{1}{2}}$$

and the bounds follow from *ii*).

It remains then to prove *i*) and *ii*).

Inequalities *i)* and *ii)* follow by applying a **Rellich-Pohozaev** identity for harmonic functions v and Lipschitz vector fields F :

$$\int_{\partial\Omega} \frac{\partial v}{\partial \nu} F \cdot \nabla v d\sigma - \frac{1}{2} \int_{\partial\Omega} |\nabla v|^2 F \cdot \nu d\sigma \\ + \frac{1}{2} \int_{\Omega} |\nabla v|^2 \operatorname{div} F dx - \int_{\Omega} (DF \cdot \nabla v) \cdot \nabla v dx = 0,$$

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We need a very specific F in order to obtain *i)* and *ii)*.

Inequalities *i)* and *ii)* are consequence of the following choice

$$F(x) := \begin{cases} 0, & \text{if } x \in \Omega \setminus \omega_h, \\ \nabla \eta, & \text{if } x \in \omega_h, \end{cases}$$

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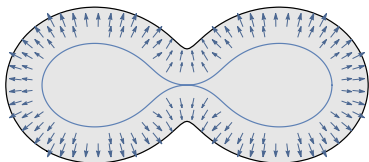
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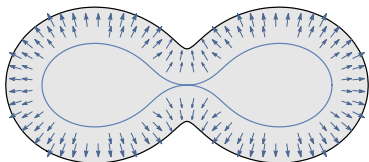
$$\omega_h := \{x \in \Omega : \text{dist}(x, \partial\Omega) < h\},$$

the number h is chosen to be the maximal possible tubular radius, and

$$\eta(x) := \frac{(h - \text{dist}(x, \partial\Omega))^2}{2}.$$



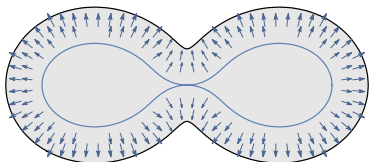
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Plugging F into the Rellich-Pohozaev identity we obtain that for harmonic functions ν it holds

$$\begin{aligned} \int_{\partial\Omega} \left(\frac{\partial \nu}{\partial \nu} \right)^2 d\sigma - \int_{\partial\Omega} |\nabla_{\partial\Omega} \nu|^2 d\sigma \\ = \frac{1}{h} \int_{\omega_h} (2(D^2 \eta \cdot \nabla \nu) \cdot \nabla \nu - |\nabla \nu|^2 \Delta \eta) dx. \end{aligned}$$



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$$\rho_i(x) = \begin{cases} \frac{(h - \text{dist}(x, \partial\Omega))\kappa_i(x')}{1 - \text{dist}(x, \partial\Omega)\kappa_i(x')}, & \text{if } i = 1, \dots, N - 1, \\ 1, & \text{if } i = N, \end{cases}$$

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where x' is the nearest point to x on $\partial\Omega$.

$$\left| \int_{\omega_h} 2(D^2\eta \cdot \nabla v) \cdot \nabla v - |\nabla v|^2 \Delta \eta dx \right| \leq (1 + (N - 1)\bar{H}_\infty h) \int_{\Omega} |\nabla v|^2 dx.$$

This and suitable estimates imply *i)* and *ii)*.

Example 1: convex domains

If Ω is a bounded and convex domain of class C^2 in \mathbb{R}^N then

$$\lambda_j \leq \sigma_j^2 + (N-1)K_\infty \sigma_j$$

and

$$\sigma_j \leq \frac{K_\infty}{2} + \sqrt{\frac{K_\infty^2}{4} + \lambda_j},$$

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$$\sigma_j \leq K_\infty + C_N \left(\frac{j}{|\partial\Omega|} \right)^{\frac{1}{N-1}}.$$

Example 2: balls

It is known that if Ω is a ball of radius R in \mathbb{R}^N , then given a Steklov eigenvalue σ it is of the form

$$\sigma = \frac{l}{R},$$

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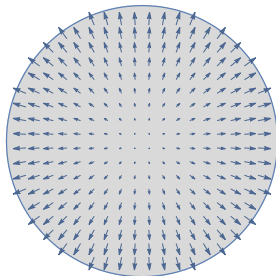
for some $l \in \mathbb{N}$. The corresponding eigenfunctions are the restriction to Ω of the harmonic polynomials in \mathbb{R}^N . Given a Laplace-Beltrami eigenvalue λ , it is of the form

$$\lambda = \frac{l(l + N - 2)}{R^2}$$

for some $l \in \mathbb{N}$. The corresponding eigenfunctions are the spherical harmonics on $\partial\Omega$.

Example 2: balls

If Ω is a ball, $\eta(x) = \frac{|x|^2}{2}$ so we can take $F(x) = x$ in the Rellich-Pohozaev identity



Example 2: balls

We obtain that for a function v harmonic in Ω

$$\int_{\partial\Omega} |\nabla_{\partial\Omega} v|^2 d\sigma = \int_{\partial\Omega} \left(\frac{\partial v}{\partial \nu} \right)^2 d\sigma + \frac{N-2}{R} \int_{\Omega} |\nabla v|^2 d\sigma$$

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- Remove the hypothesis of connected boundary;
- Less regular domains (piecewise C^1 , polygonal) with other choices of F in the Rellich-Pohozaev identity.

THANK YOU